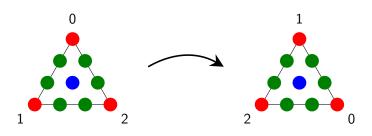
Symmetry in Finite Element Exterior Calculus

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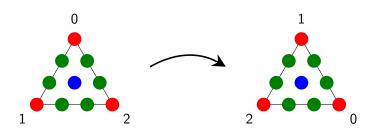
April 8-9, 2022

Symmetry of Scalar Elements



$$\mathcal{P}_3\Lambda^0(\mathcal{T}^2) = \left\langle \lambda_0^3, \lambda_1^3, \lambda_2^3, \lambda_1^2\lambda_2, \lambda_2^2\lambda_1, \lambda_2^2\lambda_0, \lambda_0^2\lambda_2, \lambda_0^2\lambda_1, \lambda_1^2\lambda_0, \lambda_0\lambda_1\lambda_2 \right\rangle.$$

Symmetry of Scalar Elements

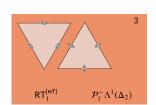


$$\mathcal{P}_3\Lambda^0(\mathcal{T}^2) = \left\langle \lambda_0^3, \lambda_1^3, \lambda_2^3, \lambda_1^2\lambda_2, \lambda_2^2\lambda_1, \lambda_2^2\lambda_0, \lambda_0^2\lambda_2, \lambda_0^2\lambda_1, \lambda_1^2\lambda_0, \lambda_0\lambda_1\lambda_2 \right\rangle.$$

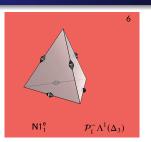
• When computing matrix of, e.g., $a(u, v) = \int_{T^2} \nabla u \cdot \nabla v$, can exploit sixfold symmetry of T^2 to compute fewer entries.

$$a\left(\lambda_0^3, \lambda_1^2 \lambda_2\right) = a\left(\lambda_0^3, \lambda_2^2 \lambda_1\right)$$
$$= a\left(\lambda_1^3, \lambda_2^2 \lambda_0\right) = a\left(\lambda_1^3, \lambda_0^2 \lambda_2\right)$$
$$= a\left(\lambda_2^3, \lambda_0^2 \lambda_1\right) = a\left(\lambda_2^3, \lambda_1^2 \lambda_0\right)$$

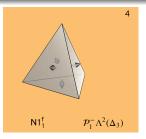
Symmetry of Vector Elements Whitney Elements



$$\langle \lambda_1 d\lambda_2 - \lambda_2 d\lambda_1, \lambda_2 d\lambda_0 - \lambda_0 d\lambda_2, \lambda_0 d\lambda_1 - \lambda_1 d\lambda_0 \rangle.$$



$$\langle \lambda_1 d\lambda_2 - \lambda_2 d\lambda_1, \\ \lambda_2 d\lambda_0 - \lambda_0 d\lambda_2, \\ \lambda_0 d\lambda_1 - \lambda_1 d\lambda_0, \\ \lambda_0 d\lambda_3 - \lambda_3 d\lambda_0, \\ \lambda_1 d\lambda_3 - \lambda_3 d\lambda_1, \\ \lambda_2 d\lambda_3 - \lambda_3 d\lambda_2 \rangle.$$

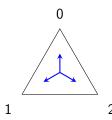


$$\langle \lambda_1 d\lambda_2 \wedge d\lambda_3 + \lambda_2 d\lambda_3 \wedge d\lambda_1 + \lambda_3 d\lambda_1 \wedge d\lambda_2, \\ \dots, \\ \lambda_0 d\lambda_1 \wedge d\lambda_2 + \lambda_1 d\lambda_2 \wedge d\lambda_0 \\ + \lambda_2 d\lambda_0 \wedge d\lambda_1 \rangle$$

Geometric symmetry \Rightarrow basis symmetry (up to sign).

Symmetry of Vector Elements

Lack of Symmetric Bases



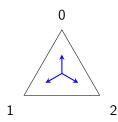
$$\mathcal{P}_0 \Lambda^1(T^2)$$

$$= \langle d\lambda_0, d\lambda_1, d\lambda_2 \rangle,$$

$$d\lambda_0 + d\lambda_1 + d\lambda_2 = 0$$

Symmetry of Vector Elements

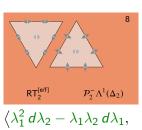
Lack of Symmetric Bases

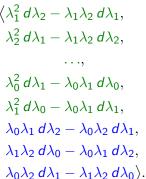


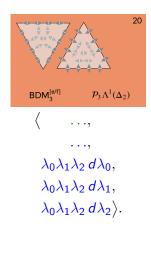
$$\mathcal{P}_0 \Lambda^1(T^2)$$

$$= \langle d\lambda_0, d\lambda_1, d\lambda_2 \rangle,$$

$$d\lambda_0 + d\lambda_1 + d\lambda_2 = 0$$







Results

Theorem (if: Licht, 2019; only if: YBK, 2021)

The following spaces have symmetry-invariant bases up to sign if and only if the corresponding condition holds.

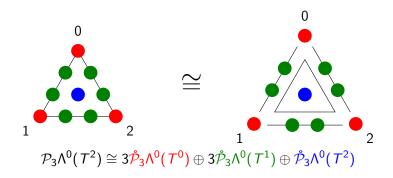
$$\mathcal{P}_r\Lambda^1(T^2)$$
 if and only if $r \notin 3\mathbb{N}_0$, $\mathcal{P}_r^-\Lambda^1(T^2)$ if and only if $r \notin 3\mathbb{N}_0 + 2$.

Theorem (YBK, 2021)

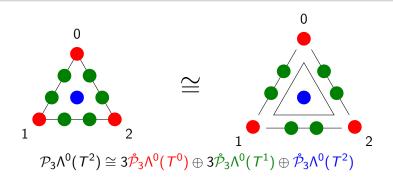
The following spaces have symmetry-invariant bases up to sign if and only if the corresponding condition holds.

$$\mathcal{P}_r \Lambda^1(T^3)$$
 always,
 $\mathcal{P}_r^- \Lambda^1(T^3)$ if and only if $r \notin 3\mathbb{N}_0 + 2$,
 $\mathcal{P}_r \Lambda^2(T^3)$ always,
 $\mathcal{P}_r^- \Lambda^2(T^3)$ always.

Methods Recursion

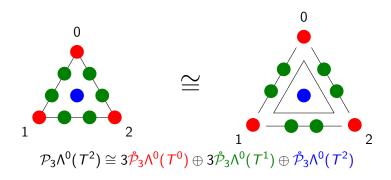


Methods Recursion



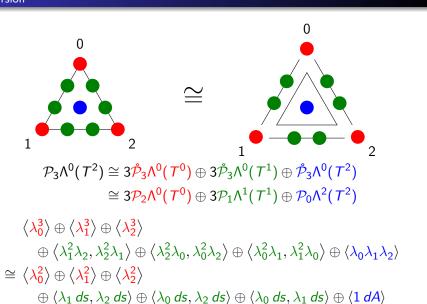
$$\begin{split} & \left< \lambda_0^3 \right> \oplus \left< \lambda_1^3 \right> \oplus \left< \lambda_2^3 \right> \\ & \oplus \left< \lambda_1^2 \lambda_2, \lambda_2^2 \lambda_1 \right> \oplus \left< \lambda_2^2 \lambda_0, \lambda_0^2 \lambda_2 \right> \oplus \left< \lambda_0^2 \lambda_1, \lambda_1^2 \lambda_0 \right> \oplus \left< \lambda_0 \lambda_1 \lambda_2 \right> \end{split}$$

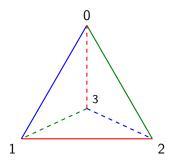
Methods Recursion

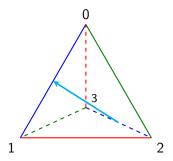


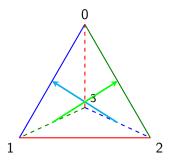
$$\begin{split} \left\langle \lambda_0^3 \right\rangle \oplus \left\langle \lambda_1^3 \right\rangle \oplus \left\langle \lambda_2^3 \right\rangle \\ & \oplus \left\langle \lambda_1^2 \lambda_2, \lambda_2^2 \lambda_1 \right\rangle \oplus \left\langle \lambda_2^2 \lambda_0, \lambda_0^2 \lambda_2 \right\rangle \oplus \left\langle \lambda_0^2 \lambda_1, \lambda_1^2 \lambda_0 \right\rangle \oplus \left\langle \lambda_0 \lambda_1 \lambda_2 \right\rangle \\ & \cong \left\langle \lambda_0^2 \right\rangle \oplus \left\langle \lambda_1^2 \right\rangle \oplus \left\langle \lambda_2^2 \right\rangle \\ & \oplus \left\langle \lambda_1 \, ds, \lambda_2 \, ds \right\rangle \oplus \left\langle \lambda_0 \, ds, \lambda_2 \, ds \right\rangle \oplus \left\langle \lambda_0 \, ds, \lambda_1 \, ds \right\rangle \oplus \left\langle 1 \, dA \right\rangle \end{split}$$

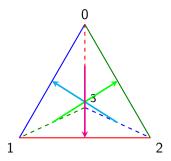
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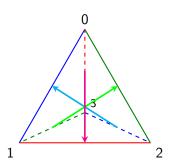








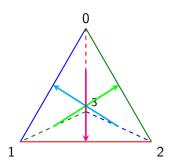
Tetrahedron Basis



$$\mathcal{P}_0 \Lambda^1 (T^3)$$

$$= \langle d\lambda_0 + d\lambda_1 - d\lambda_2 - d\lambda_3, d\lambda_0 + d\lambda_2 - d\lambda_1 - d\lambda_3, d\lambda_1 + d\lambda_2 - d\lambda_0 - d\lambda_3 \rangle$$

$$=: \langle \alpha, \beta, \gamma \rangle.$$



$$\mathcal{P}_0 \Lambda^1(T^3)$$

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$$=: \langle \alpha, \beta, \gamma \rangle.$$

$$\mathcal{P}_{2}\Lambda^{1}(T^{3})$$

$$= \mathcal{P}_{2}\Lambda^{0}(T^{3}) \otimes \mathcal{P}_{0}\Lambda^{1}(T^{3})$$

$$= \langle \lambda_{0}^{2}\alpha, \lambda_{0}^{2}\beta, \lambda_{0}^{2}\gamma, \lambda_{1}^{2}\alpha, \lambda_{1}^{2}\beta, \lambda_{1}^{2}\gamma, \lambda_{2}^{2}\alpha, \lambda_{2}^{2}\beta, \lambda_{2}^{2}\gamma, \lambda_{3}^{2}\alpha, \lambda_{3}^{2}\beta, \lambda_{3}^{2}\gamma, \lambda_{0}\lambda_{1}\alpha, \lambda_{0}\lambda_{1}\beta, \lambda_{0}\lambda_{1}\gamma, \lambda_{0}\lambda_{2}\alpha, \lambda_{0}\lambda_{2}\beta, \lambda_{0}\lambda_{2}\gamma, \lambda_{0}\lambda_{3}\alpha, \lambda_{0}\lambda_{3}\beta, \lambda_{0}\lambda_{3}\gamma, \lambda_{1}\lambda_{2}\alpha, \lambda_{1}\lambda_{2}\beta, \lambda_{1}\lambda_{2}\gamma, \lambda_{1}\lambda_{3}\alpha, \lambda_{1}\lambda_{3}\beta, \lambda_{1}\lambda_{3}\gamma, \lambda_{2}\lambda_{3}\alpha, \lambda_{2}\lambda_{3}\beta, \lambda_{2}\lambda_{3}\gamma\rangle.$$

Representations of $\mathbb{Z}/3$

- The 1D representation ${\bf 1}$ where $\mathbb{Z}/3$ acts trivially.
- The 2D representation **2** where $\mathbb{Z}/3$ acts by 120° rotations.
- The 3D representation **3** where $\mathbb{Z}/3$ acts by permuting the coordinates.
 - $\mathbf{3} \cong \mathbf{1} \oplus \mathbf{2}$ because $\langle (1,1,1) \rangle$ is an invariant subspace.

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Invariant bases

1 and 3 have symmetry-invariant bases, but 2 does not.



Proposition

A representation $V \cong m\mathbf{1} \oplus n\mathbf{2}$ has a $\mathbb{Z}/3$ -invariant basis up to sign if and only if $m \geq n$.

Thank You and References



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